

ESSENTIAL NORM ESTIMATES FOR LITTLE HANKEL OPERATORS ON

$$L_a^2(\mathbb{C}_+)^*$$

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Abstract

In this paper, we give estimates for the essential norm of a bounded little Hankel operator defined on the Bergman space of the right half plane. As an application of these estimates, we also give a necessary and sufficient condition for the little Hankel operator to be compact.

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1 Introduction

Let $\mathbb{C}_+ = \{s = x + iy \in \mathbb{C} : \text{Re } s > 0\}$ be the right half plane. Let $d\mu(s) = dx dy$ be the area measure. Let $L^2(\mathbb{C}_+, d\mu)$ be the space of complex-valued, square-integrable, measurable functions on \mathbb{C}_+ with respect to the area measure. Let $L_a^2(\mathbb{C}_+)$ be the closed subspace [1] of $L^2(\mathbb{C}_+, d\mu)$ consisting of those functions in $L^2(\mathbb{C}_+, d\mu)$ that are analytic. The space $L_a^2(\mathbb{C}_+)$ is referred to as the Bergman space of the right half plane. The functions

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$H(s, w) = \frac{1}{(s+\bar{w})^2}$, $s \in \mathbb{C}_+$, $w \in \mathbb{C}_+$ is the reproducing kernel [2] for $L_a^2(\mathbb{C}_+)$. Let $L^\infty(\mathbb{C}_+)$ be the space of complex-valued, essentially bounded, Lebesgue measurable functions on \mathbb{C}_+ . For $f \in L^\infty(\mathbb{C}_+)$, $\|f\|_\infty = \operatorname{ess\,sup}_{s \in \mathbb{C}_+} |f(s)| < \infty$.

The space $L^\infty(\mathbb{C}_+)$ is a Banach space with respect to the essential supremum norm. For $\phi \in L^\infty(\mathbb{C}_+)$, we define the multiplication operator \mathcal{M}_ϕ from $L^2(\mathbb{C}_+, d\mu)$ into $L^2(\mathbb{C}_+, d\mu)$ by $(\mathcal{M}_\phi f)(s) = \phi(s)f(s)$ and the little Hankel operator \tilde{h}_ϕ is a mapping from $L_a^2(\mathbb{C}_+)$ into $\overline{L_a^2(\mathbb{C}_+)}$ defined by $\tilde{h}_\phi f = \overline{P_+(\phi f)}$, where $\overline{P_+}$ is the projection operator from $L^2(\mathbb{C}_+, d\mu)$ onto $\overline{L_a^2(\mathbb{C}_+)} = \{\bar{f} : f \in L_a^2(\mathbb{C}_+)\}$. There are also many equivalent ways of defining little Hankel operators on $L_a^2(\mathbb{C}_+)$. Let \mathcal{S}_ϕ be the mapping from $L_a^2(\mathbb{C}_+)$ into $L_a^2(\mathbb{C}_+)$ defined by $\mathcal{S}_\phi f = P_+(\mathcal{J}(\phi f))$, where P_+ denote the orthogonal projection from $L^2(\mathbb{C}_+, d\mu)$ onto $L_a^2(\mathbb{C}_+)$ and \mathcal{J} is the mapping from $L^2(\mathbb{C}_+, d\mu)$ into $L^2(\mathbb{C}_+, d\mu)$ such that $\mathcal{J}f(s) = f(\bar{s})$. Notice that \mathcal{J} is unitary and $\mathcal{J}\mathcal{S}_\phi f = \mathcal{J}(P_+(\mathcal{J}(\phi f))) = \mathcal{J}P_+\mathcal{J}(\phi f) = \overline{P_+(\phi f)} = \tilde{h}_\phi f$ for $f \in L_a^2(\mathbb{C}_+)$. Let Γ_ϕ be the mapping from $L_a^2(\mathbb{C}_+)$ into $L_a^2(\mathbb{C}_+)$ defined by $\Gamma_\phi f = P_+\mathcal{M}_\phi\mathcal{J}f$. Thus $\Gamma_\phi f = P_+\mathcal{M}_\phi\mathcal{J}f = P_+(\phi(s)f(\bar{s})) = P_+(\mathcal{J}(\phi(\bar{s})f(s))) = \mathcal{S}_{\mathcal{J}\phi}f$ for all $f \in L_a^2(\mathbb{C}_+)$. Hence $\Gamma_\phi f = \mathcal{S}_{\mathcal{J}\phi}f$. Thus we obtain $\tilde{h}_\phi = \mathcal{J}\mathcal{S}_\phi$ and $\Gamma_\phi = \mathcal{S}_{\mathcal{J}\phi}$. Since \mathcal{J} is unitary, the three operators \tilde{h}_ϕ , \mathcal{S}_ϕ and Γ_ϕ are referred to as little Hankel operators on $L_a^2(\mathbb{C}_+)$ and a given result on little Hankel operators can be stated using the operators \tilde{h}_ϕ , \mathcal{S}_ϕ and Γ_ϕ .

Let $\mathbb{D} = \{z \in \mathbb{C} : |z| < 1\}$ be the open unit disk in the complex plane \mathbb{C} . Let $L^2(\mathbb{D}, dA)$ be the space of complex-valued, square-integrable, measurable functions on \mathbb{D} with respect to the normalized area measure $dA(z) = \frac{1}{\pi}dxdy$. Let $L_a^2(\mathbb{D})$ be the space consisting of those functions of $L^2(\mathbb{D}, dA)$ that are analytic. The space $L_a^2(\mathbb{D})$ is a closed subspace of $L^2(\mathbb{D}, dA)$ and is called the Bergman space of the open unit disk \mathbb{D} . Let $L^\infty(\mathbb{D})$ be the space of complex-valued, essentially bounded, Lebesgue measurable functions on \mathbb{D} with the essential supremum norm. For $\phi \in L^\infty(\mathbb{D})$, the multiplication operator M_ϕ from $L^2(\mathbb{D}, dA)$ into $L^2(\mathbb{D}, dA)$ is defined by $M_\phi f = \phi f$ and the little Hankel operator h_ϕ is a mapping from $L_a^2(\mathbb{D})$ into $\overline{L_a^2(\mathbb{D})}$ defined by $h_\phi f = \overline{P(\phi f)}$, where \overline{P} is the projection operator from $L^2(\mathbb{D}, dA)$ onto $\overline{L_a^2(\mathbb{D})} = \{\bar{f} : f \in L_a^2(\mathbb{D})\}$. Let S_ϕ be the mapping from $L_a^2(\mathbb{D})$ into $L_a^2(\mathbb{D})$ defined by $S_\phi f = P(J(\phi f))$, where P is the orthogonal projection from $L^2(\mathbb{D}, dA)$ onto $L_a^2(\mathbb{D})$ and J is the mapping from $L^2(\mathbb{D}, dA)$ into itself such that $Jf(z) = f(\bar{z})$. Notice that J is unitary and $JS_\phi f = J(P(J(\phi f))) = JPJ(\phi f) = \overline{P(\phi f)} = h_\phi f$ for all $f \in L_a^2(\mathbb{D})$. Let Γ_ϕ be the mapping from $L_a^2(\mathbb{D})$ into $L_a^2(\mathbb{D})$ defined by $\Gamma_\phi f = PM_\phi Jf$, where M_ϕ

is the mapping from $L^2(\mathbb{D}, dA)$ into $L^2(\mathbb{D}, dA)$ defined by $M_\phi f = \phi f$. Thus $\Gamma_\phi f = PM_\phi Jf = P(\phi(z)f(\bar{z})) = P(J(\phi(\bar{z})f(z))) = S_{J\phi}f$ for all $f \in L^2_a(\mathbb{D})$. Hence $\Gamma_\phi = S_{J\phi}$. Since J is unitary, the three operators h_ϕ, S_ϕ and Γ_ϕ are referred to as little Hankel operators on $L^2_a(\mathbb{D})$. The sequence of functions $\{e_n(z)\}_{n=0}^\infty = \{\sqrt{n+1}z^n\}_{n=0}^\infty$ form an orthonormal basis for $L^2_a(\mathbb{D})$. Since point evaluation at $z \in \mathbb{D}$ is a bounded linear functional on the Hilbert space $L^2_a(\mathbb{D})$, the Riesz representation theorem implies that there exists a unique function K_z in $L^2_a(\mathbb{D})$ such that

$$f(z) = \int_{\mathbb{D}} f(w)\overline{K_z(w)}dA(w).$$

for all f in $L^2_a(\mathbb{D})$. Let $K(z, w)$ be the function on $\mathbb{D} \times \mathbb{D}$ defined by

$$K(z, w) = \overline{K_z(w)}.$$

The function $K(z, w)$ is analytic in z and co-analytic in w . Since

$$f(z) = \int_{\mathbb{D}} f(w)K(z, w)dA(w), f \in L^2_a(\mathbb{D}),$$

the function $K(z, w) = \frac{1}{(1-z\bar{w})^2}$, $z, w \in \mathbb{D}$ and is the reproducing kernel [7] of $L^2_a(\mathbb{D})$. For $a \in \mathbb{D}$, let $k_a(z) = \frac{K(z, a)}{\sqrt{K(a, a)}} = \frac{(1-|a|^2)}{(1-\bar{a}z)^2}$. The function k_a is called the normalized reproducing kernel for $L^2_a(\mathbb{D})$. It is clear that $\|k_a\|_2 = 1$. Let $Aut(\mathbb{D})$ be the Lie group of all automorphisms (biholomorphic mappings) of \mathbb{D} . We can define for each $a \in \mathbb{D}$ an automorphism ϕ_a in $Aut(\mathbb{D})$ such that

- (i) $(\phi_a \circ \phi_a)(z) = z$;
- (ii) $\phi_a(0) = a, \phi_a(a) = 0$;
- (iii) ϕ_a has a unique fixed point in \mathbb{D} .

In fact, $\phi_a(z) = \frac{a-z}{1-\bar{a}z}$ for all a and z in \mathbb{D} . An easy calculation shows that the derivative of ϕ_a at z is equal to $-k_a(z)$. It follows that the real Jacobian determinant of ϕ_a at z is $J_{\phi_a}(z) = |k_a(z)|^2 = \frac{(1-|a|^2)^2}{|1-\bar{a}z|^4}$. Given $a \in \mathbb{D}$ and f any measurable function on \mathbb{D} , we define a function $U_a f$ on \mathbb{D} by $U_a f(z) = k_a(z)f(\phi_a(z))$. Let $\mathcal{L}(H)$ be the set of all bounded linear operators from the Hilbert space H into itself. Let $\mathcal{LC}(H)$ be the set of all compact operators in $\mathcal{L}(H)$. The essential norm of an operator $T \in \mathcal{L}(H)$ is the distance of the operator from the space of compact operators, that is

$$\|T\|_e = \inf\{\|T - K\| : K \text{ is compact.}\}$$

In particular, T is compact if and only if $\|T\|_e = 0$. Essential norm estimates for bounded operators on the Bergman space are studied in [4] and [6]. The layout of this paper is as follows. In §2, we introduce a class of unitary operators defined on $L_a^2(\mathbb{C}_+)$ induced by the automorphisms $t_a(s)$ of \mathbb{C}_+ . In §3, we introduce the functions $B(s, w), B_{\bar{w}}(s)$ and $b_{\bar{w}}(s), s, w \in \mathbb{C}_+$ and establish relations between them. We also show that the function $B(s, w)$ satisfy an inequality like the Bergman kernel (see [3]) $K(z, w)$ defined for the space $L_a^2(\mathbb{D})$. In §4, we introduce the operators Q_1 and \mathcal{V}_1 and show that they are bounded on $L^2(\mathbb{C}_+, d\mu)$. In §5, we establish that if $\phi \in L^2(\mathbb{C}_+, d\mu)$, then the little Hankel operator $h_{\bar{\phi}}$ is bounded if and only if $\mathcal{V}_1\phi$ is bounded on \mathbb{C}_+ . In §6, we give estimates for the essential norm of bounded little Hankel operators on the Bergman space $L_a^2(\mathbb{C}_+, d\mu)$ in terms of the function $\mathcal{V}_1\phi$ and applications of the result are also obtained.

2 A class of unitary operators on $L_a^2(\mathbb{C}_+)$

In this section, we introduce a class of unitary operators defined on $L_a^2(\mathbb{C}_+)$ induced by the automorphisms $t_a(s)$ of \mathbb{C}_+ .

Define $M : \mathbb{C}_+ \rightarrow \mathbb{D}$ by $Ms = \frac{1-s}{1+s}$. Then M is one-one, onto and $M^{-1} : \mathbb{D} \rightarrow \mathbb{C}_+$ is given by $M^{-1}(z) = \frac{1-z}{1+z}$. Thus M is its self-inverse. Let $W : L_a^2(\mathbb{D}) \rightarrow L_a^2(\mathbb{C}_+)$ be defined by $Wg(s) = \frac{2}{\sqrt{\pi}}g(Ms)\frac{1}{(1+s)^2}$. The map W is one-one and onto. Hence W^{-1} exists and $W^{-1} : L_a^2(\mathbb{C}_+) \rightarrow L_a^2(\mathbb{D})$ is given by $W^{-1}G(z) = 2\sqrt{\pi}G(Mz)\frac{1}{(1+z)^2}$, where $Mz = \frac{1-z}{1+z}$.

Lemma 1. *If $a \in \mathbb{D}$ and $a = c + id, c, d \in \mathbb{R}$, then the following hold:*

- (i) $t_a(s) = \frac{-ids+(1-c)}{(1+c)s+id}$ is an automorphism from \mathbb{C}_+ onto \mathbb{C}_+ .
- (ii) $(t_a \circ t_a)(s) = s$.
- (iii) $t'_a(s) = -l_a(s)$, where $l_a(s) = \frac{1-|a|^2}{((1+c)s+id)^2}$.

Proof. This can be verified by direct calculations. □

For $a \in \mathbb{D}$, define $V_a : L_a^2(\mathbb{C}_+) \rightarrow L_a^2(\mathbb{C}_+)$ by $(V_a g)(s) = (g \circ t_a)(s)l_a(s)$. In Proposition 1, we show that V_a is a self-adjoint, unitary operator which is also an involution.

Proposition 1. *For $a \in \mathbb{D}$,*

- (i) $V_a l_a = 1$.

(ii) $V_a^{-1} = V_a$ and V_a is an involution, i.e. $V_a^2 = I_{\mathcal{L}(L_a^2(\mathbb{C}_+))}$, where $I_{\mathcal{L}(L_a^2(\mathbb{C}_+)}$ is the identity operator from $L_a^2(\mathbb{C}_+)$ into itself.

(iii) V_a is self-adjoint.

(iv) V_a is unitary, $\|V_a\| = 1$.

(v) $V_a P_+ = P_+ V_a$.

Proof. One can prove (i), (ii), (iii) and (iv) by direct calculations. Notice that V_a can also be defined from $L^2(\mathbb{C}_+)$ into itself. To prove (v), observe that $V_a(L_a^2(\mathbb{C}_+)) \subset L_a^2(\mathbb{C}_+)$ and $V_a(L_a^2(\mathbb{C}_+))^\perp \subset (L_a^2(\mathbb{C}_+))^\perp$. Now let $f \in L^2(\mathbb{C}_+)$ and $f = f_1 + f_2$, where $f_1 \in L_a^2(\mathbb{C}_+)$ and $f_2 \in (L_a^2(\mathbb{C}_+))^\perp$. Hence,

$$P_+ V_a f = P_+ V_a (f_1 + f_2) = P_+ (V_a f_1 + V_a f_2) = P_+ V_a f_1 = V_a f_1 = V_a P_+ f.$$

□

3 The function $B(s, w)$

In this section, we introduce the functions $B(s, w)$ and $b_{\bar{w}}(s)$, $s, w \in \mathbb{C}_+$ and establish relations between them. We also show that the function $B(s, w)$ satisfy an inequality like the Bergman kernel (see [3]) $K(z, w)$ defined for the space $L_a^2(\mathbb{D})$.

Suppose $a \in \mathbb{D}$ and $w = \frac{1-\bar{a}}{1+a} = M\bar{a} \in \mathbb{C}_+$. Define $b_{\bar{w}}(s) = \frac{1}{\sqrt{\pi}} \frac{1+w}{1+\bar{w}} \frac{2Re w}{(s+w)^2}$. Let $B(s, w) = B_{\bar{w}}(s) = \frac{1}{\pi} \frac{(1+a)^2}{(1-\bar{a}Ms)^2} \frac{1}{(1+s)^2}$.

Lemma 2. *Let $s, w \in \mathbb{C}_+$. The following hold:*

(i) $(b_{\bar{w}}(\bar{w}))^2 = B(\bar{w}, w)$.

(ii) $|b_{\bar{w}}(s)| \|B_{\bar{w}}\| = |B_{\bar{w}}(s)|$.

Proof. Let $w \in \mathbb{C}_+$ and $w = M\bar{a} = \frac{1-\bar{a}}{1+\bar{a}}$. Since

$$\begin{aligned}
 b_{\bar{w}}(s) &= \frac{1}{\sqrt{\pi}} \frac{1+w}{1+\bar{w}} \frac{2\operatorname{Re}w}{[s+w]^2} = \frac{2}{\sqrt{\pi}} \frac{\operatorname{Re}w}{(1+w)(1+\bar{w})} \frac{(1+w)^2}{[s+w]^2} \\
 &= \frac{2}{\sqrt{\pi}} \frac{4\operatorname{Re}w}{|1+w|^2} \frac{(1+w)^2}{4} \frac{1}{[s+w]^2} = \frac{2}{\sqrt{\pi}} \frac{\frac{|1+w|^2 - |1-w|^2}{|1+w|^2}}{\left[\frac{2}{(1+w)}\right]^2} \frac{1}{[s+w]^2} \\
 &= \frac{2}{\sqrt{\pi}} \frac{1 - \left|\frac{1-w}{1+w}\right|^2}{\left(1 + \frac{1-w}{1+w}\right)^2} \frac{1}{[s+w]^2} = \frac{2}{\sqrt{\pi}} \frac{1 - |a|^2}{(1+\bar{a})^2} \frac{1}{[s+w]^2}, \text{ where } \frac{1-\bar{a}}{1+\bar{a}} = w, \\
 &= \frac{2}{\sqrt{\pi}} \frac{1 - |a|^2}{(1+\bar{a})^2} \frac{1}{\left[s + \frac{1-\bar{a}}{1+\bar{a}}\right]^2} = \frac{2}{\sqrt{\pi}} \frac{1 - |a|^2}{(1+\bar{a})^2 \left[s + \frac{1-\bar{a}}{1+\bar{a}}\right]^2} \\
 &= \frac{2}{\sqrt{\pi}} \frac{1 - |a|^2}{[1 - \bar{a} + s(1+\bar{a})]^2} = \frac{2}{\sqrt{\pi}} \frac{1 - |a|^2}{[1 + s - \bar{a} + \bar{a}s]^2} \\
 &= \frac{(-1)}{\sqrt{\pi}} \frac{(1 - |a|^2)(1+s)^2}{[1 + s - \bar{a} + \bar{a}s]^2} \frac{(-2)}{(1+s)^2} = \frac{(-1)}{\sqrt{\pi}} \frac{1 - |a|^2}{\left[1 - \bar{a}\frac{1-s}{1+s}\right]^2} \frac{(-2)}{(1+s)^2} \\
 &= \frac{2}{\sqrt{\pi}} \frac{1 - |a|^2}{[1 - \bar{a}(Ms)]^2} \frac{1}{(1+s)^2},
 \end{aligned}$$

we obtain

$$\begin{aligned}
 b_{\bar{w}}(\bar{w}) &= \frac{2}{\sqrt{\pi}} \frac{(1 - |a|^2)}{(1 - \bar{a}M\bar{w})^2} \frac{1}{(1 + \bar{w})^2} = \frac{2}{\sqrt{\pi}} \frac{(1 - |a|^2)}{(1 - |a|^2)^2} \frac{1}{\left(1 + \frac{1-\bar{a}}{1+\bar{a}}\right)^2} \\
 &= \frac{2}{\sqrt{\pi}} \frac{1}{(1 - |a|^2)} \frac{(1+a)^2}{4} = \frac{1}{2\sqrt{\pi}} \frac{(1+a)^2}{(1 - |a|^2)}.
 \end{aligned}$$

Thus

$$\begin{aligned}
 b_{\bar{w}}(s)b_{\bar{w}}(\bar{w}) &= \frac{2}{\sqrt{\pi}} \frac{(1 - |a|^2)}{(1 - \bar{a}Ms)^2} \frac{1}{(1+s)^2} \frac{1}{2\sqrt{\pi}} \frac{(1+a)^2}{(1 - |a|^2)} \\
 &= \frac{1}{\pi} \frac{1}{(1 - \bar{a}Ms)^2} \frac{(1+a)^2}{(1+s)^2} = \frac{(-1)}{2\pi} \frac{(1+a)^2}{(1 - \bar{a}Ms)^2} \frac{(-2)}{(1+s)^2} \\
 &= \frac{(-1)}{2\pi} \frac{(1+a)^2}{(1 - \bar{a}Ms)^2} M' = B(s, w).
 \end{aligned}$$

Thus $b_{\bar{w}}(s) = \frac{B(s, w)}{b_{\bar{w}}(\bar{w})}$ and $(b_{\bar{w}}(\bar{w}))^2 = B(\bar{w}, w)$. This proves (i). To prove (ii),

notice that

$$\begin{aligned} \|B_{\bar{w}}\|^2 &= \langle B_{\bar{w}}, B_{\bar{w}} \rangle = \int_{\mathbb{C}_+} |B_{\bar{w}}(s)|^2 d\mu(s) = \int_{\mathbb{C}_+} |B(s, w)|^2 d\mu(s) \\ &= \int_{\mathbb{C}_+} |b_{\bar{w}}(\bar{w})|^2 |b_{\bar{w}}(s)|^2 d\mu(s) = |b_{\bar{w}}(\bar{w})|^2 \int_{\mathbb{C}_+} |b_{\bar{w}}(s)|^2 d\mu(s) \\ &= |b_{\bar{w}}(\bar{w})|^2 \|b_{\bar{w}}\|_2^2 = |b_{\bar{w}}(\bar{w})|^2, \end{aligned}$$

since $\|b_{\bar{w}}\|_2 = 1$. Thus $\|B_{\bar{w}}\| = |b_{\bar{w}}(\bar{w})|$ and hence $|b_{\bar{w}}(s)| \|B_{\bar{w}}\| = |B_{\bar{w}}(s)|$. \square

Lemma 3. *Suppose $-\frac{1}{2} < q < p - 1$. Then there exists a positive constant C such that*

$$\int_{\mathbb{C}_+} |B(\bar{s}, \bar{w})|^p |B(\bar{w}, w)|^{-q} d\mu(\bar{w}) \leq C |B(\bar{s}, s)|^{p-q-1}$$

for all $s \in \mathbb{C}_+$.

Proof. Since $B(s, w) = \frac{1}{\pi} \frac{(1+a)^2}{(1-\bar{a}Ms)^2} \frac{1}{(1+s)^2}$ and $Ma = \bar{w}$, we obtain

$$\begin{aligned} &\int_{\mathbb{C}_+} |B(\bar{s}, \bar{w})|^p |B(\bar{w}, w)|^{-q} d\mu(\bar{w}) \\ &= \int_{\mathbb{C}_+} \left| \frac{1}{\pi} \frac{(1+\bar{a})^2}{(1-aM\bar{s})^2} \frac{1}{(1+\bar{s})^2} \right|^p \left| \frac{1}{4\pi} \frac{(1+a)^4}{(1-|a|^2)^2} \right|^{-q} d\mu(Ma) \\ &= \int_{\mathbb{D}} \left| \frac{1}{\pi} \frac{(1+\bar{a})^2}{(1-a\bar{z})^2} \frac{1}{(1+M\bar{z})^2} \right|^p \left| \frac{1}{4\pi} \frac{(1+a)^4}{(1-|a|^2)^2} \right|^{-q} \left| \frac{(-2)}{(1+z)^2} \right|^2 dA(a) \\ &= \int_{\mathbb{D}} \left| \frac{1}{\pi} \frac{(1+\bar{a})^2}{(1-a\bar{z})^2} \frac{1}{\left(1 + \frac{1-\bar{z}}{1+\bar{z}}\right)^2} \right|^p \left| \frac{1}{4\pi} \frac{(1+a)^4}{(1-|a|^2)^2} \right|^{-q} \left| \frac{(-2)}{(1+a)^2} \right|^2 dA(a) \\ &= \int_{\mathbb{D}} \left| \frac{1}{\pi} \frac{(1+\bar{a})^2}{(1-a\bar{z})^2} \frac{(1+\bar{z})^2}{(1+\bar{z}+1-\bar{z})^2} \right|^p \left| \frac{1}{4\pi} \frac{(1+a)^4}{(1-|a|^2)^2} \right|^{-q} \left| \frac{(-2)}{(1+a)^2} \right|^2 dA(a) \\ &= \int_{\mathbb{D}} \left| \frac{1}{4\pi} \frac{(1+\bar{a})^2(1+\bar{z})^2}{(1-a\bar{z})^2} \right|^p \left| \frac{1}{4\pi} \frac{(1+a)^4}{(1-|a|^2)^2} \right|^{-q} \left| \frac{(-2)}{(1+a)^2} \right|^2 dA(a) \end{aligned}$$

$$\begin{aligned}
&= \left(\frac{1}{4\pi}\right)^{p-q} 2^2 |1+\bar{z}|^{2p} \int_{\mathbb{D}} \left| \frac{1}{(1-\bar{a}z)^2} \right|^p \left| \frac{1}{(1-|a|^2)^2} \right|^{-q} |1+a|^{-4q-4} |1+\bar{a}|^{2p} dA(a) \\
&= \left(\frac{1}{4\pi}\right)^{p-q} 2^2 |1+\bar{z}|^{2p} \int_{\mathbb{D}} |K(z, a)|^p |K(a, a)|^{-q} |1+a|^{-4q-4} |1+\bar{a}|^{2p} dA(a) \\
&\leq \frac{4}{(4\pi)^{p-q}} 2^{2p} 2^{2p} 2^{-4q-4} \int_{\mathbb{D}} |K(z, a)|^p |K(a, a)|^{-q} dA(a) \\
&\leq \frac{1}{4} \left(\frac{4}{\pi}\right)^{p-q} \int_{\mathbb{D}} |K(z, a)|^p |K(a, a)|^{-q} dA(a).
\end{aligned}$$

From [3], we obtain

$$\int_{\mathbb{C}_+} |B(\bar{s}, \bar{w})|^p |B(\bar{w}, w)|^{-q} d\mu(\bar{w}) \leq C \frac{1}{4} \left(\frac{4}{\pi}\right)^{p-q} K(z, z)^{p-q-1}$$

for some constant C . Let $C_1 = C \frac{1}{4} \left(\frac{4}{\pi}\right)^{p-q}$. Then

$$\begin{aligned}
&\int_{\mathbb{C}_+} |B(\bar{s}, \bar{w})|^p |B(\bar{w}, w)|^{-q} d\mu(\bar{w}) \leq C_1 K(z, z)^{p-q-1} = C_1 K_z(z)^{p-q-1} \\
&= C_1 \langle K_z, K_z \rangle^{p-q-1} = C_1 \|K_z\|^{2(p-q-1)} \left\langle \frac{K_z}{\|K_z\|}, \frac{K_z}{\|K_z\|} \right\rangle^{p-q-1} \\
&= C_2 \langle k_z, k_z \rangle^{p-q-1} \text{ where } C_2 = C_1 \|K_z\|^{2(p-q-1)}.
\end{aligned}$$

Thus, if $z = M\bar{s}$, then

$$\begin{aligned}
\int_{\mathbb{C}_+} |B(\bar{s}, \bar{w})|^p |B(\bar{w}, w)|^{-q} d\mu(\bar{w}) &\leq C_2 \left(\frac{\|K_z\|^2}{\|B_{\bar{s}}\|^2} |B(\bar{s}, s)| \right)^{p-q-1} \\
&= C_3 |B(\bar{s}, s)|^{p-q-1},
\end{aligned}$$

where $C_3 = C_2 \frac{\|K_z\|^{2(p-q-1)}}{\|B_{\bar{s}}\|^{2(p-q-1)}}$. This complete the proof. \square

Lemma 4. Let $s, w \in \mathbb{C}_+$, and $w = M\bar{a}$. Then $|B(\bar{s}, \bar{w})| = |B(\bar{w}, \bar{s})|$.

Proof. Let $s, w \in \mathbb{C}_+$ and $w = M\bar{a}$. Since $B(s, w) = \frac{1}{\pi} \frac{1}{(1-\bar{a}Ms)^2} \frac{(1+a)^2}{(1+s)^2}$, we

obtain

$$\begin{aligned}
B(\bar{s}, \bar{w}) &= \frac{1}{\pi} \frac{1}{(1 - aM\bar{s})^2} \frac{(1 + Mw)^2}{(1 + \bar{s})^2} = \frac{1}{\pi} \frac{\frac{(1+w+1-w)^2}{(1+w)^2}}{\left(1 - a\frac{1-\bar{s}}{1+\bar{s}}\right)^2} \frac{1}{(1 + \bar{s})^2} \\
&= \frac{4}{\pi} \frac{1}{(1 + w)^2} \frac{(1 + \bar{s})^2}{(1 + \bar{s} - a + a\bar{s})^2} \frac{1}{(1 + \bar{s})^2} = \frac{4}{\pi} \frac{1}{(1 + w)^2} \frac{1}{(1 - a + \bar{s}(1 + a))^2} \\
&= \frac{4}{\pi} \frac{1}{(1 + w)^2} \frac{1}{(1 + a)^2 \left(\frac{1-a}{1+a} + \bar{s}\right)^2} = \frac{4}{\pi} \frac{1}{(1 + w)^2} \frac{1}{(1 + a)^2 (\bar{s} + \bar{w})^2} \\
&= \frac{4}{\pi} \frac{1}{(1 + w)^2} \frac{1}{(1 + M\bar{w})^2 (\bar{s} + \bar{w})^2} = \frac{4}{\pi} \frac{(1 + \bar{w})^2}{4(1 + w)^2} \frac{1}{(\bar{s} + \bar{w})^2} \\
&= \frac{1}{\pi} \left(\frac{1 + \bar{w}}{1 + w}\right)^2 \frac{1}{(\bar{s} + \bar{w})^2}.
\end{aligned}$$

Similarly,

$$\begin{aligned}
B(\bar{w}, \bar{s}) &= \frac{4}{\pi} \frac{1}{(1 + s)^2} \frac{1}{(1 + M\bar{s})^2 (\bar{s} + \bar{w})^2} = \frac{4}{\pi} \frac{1}{(1 + s)^2} \frac{1}{(\bar{s} + \bar{w})^2} \frac{1}{\left(1 + \frac{1-\bar{s}}{1+\bar{s}}\right)^2} \\
&= \frac{4}{\pi} \frac{1}{(1 + s)^2} \frac{1}{(\bar{s} + \bar{w})^2} \frac{(1 + \bar{s})^2}{4} = \frac{1}{\pi} \frac{(1 + \bar{s})^2}{(1 + s)^2} \frac{1}{(\bar{s} + \bar{w})^2}.
\end{aligned}$$

Thus $|B(\bar{s}, \bar{w})| = |B(\bar{w}, \bar{s})|$ for all $s, w \in \mathbb{C}_+$.

□

4 Integral operator

In this section, we introduce the operators Q_1 and \mathcal{V}_1 and prove that these operators are bounded on $L^2(\mathbb{C}_+, d\mu)$. For $\phi \in L^2(\mathbb{C}_+, d\mu)$ and $s \in \mathbb{C}_+$, we define

$$Q_1\phi(s) = 3 \int_{\mathbb{C}_+} \frac{|B(\bar{s}, \bar{w})|^2}{|B(\bar{w}, w)|} \phi(w) d\mu(w)$$

and

$$\mathcal{V}_1\phi(w) = 3 \int_{\mathbb{C}_+} \frac{|B(\bar{s}, \bar{w})|^2}{|B(\bar{w}, w)|} \phi(s) d\mu(s).$$

Proposition 2. *The operators Q_1 and \mathcal{V}_1 are bounded on $L^2(\mathbb{C}_+, d\mu)$.*

Proof. Notice that the boundedness of Q_1 follows from the boundedness of \mathcal{V}_1 . Thus we only show that Q_1 is a bounded operator on $L^2(\mathbb{C}_+, d\mu)$. Take $t > 0$ and let $h(s) = |B(\bar{s}, s)|^t$. Then by Lemma 3 and Lemma 4, we obtain

$$\begin{aligned} \int_{\mathbb{C}_+} \frac{|B(\bar{s}, \bar{w})|^2}{|B(\bar{w}, w)|} h(s) d\mu(s) &= |B(\bar{w}, w)|^{-1} \int_{\mathbb{C}_+} |B(\bar{s}, \bar{w})|^2 |B(\bar{s}, s)|^t d\mu(s) \\ &= |B(\bar{w}, w)|^{-1} \int_{\mathbb{C}_+} |B(\bar{w}, \bar{s})|^2 |B(\bar{s}, s)|^t d\mu(s) \\ &\leq |B(\bar{w}, w)|^{-1} C |B(\bar{w}, w)|^{2+t-1} \\ &= C |B(\bar{w}, w)|^t = Ch(w); \end{aligned} \quad (1)$$

and

$$\begin{aligned} \int_{\mathbb{C}_+} \frac{|B(\bar{s}, \bar{w})|^2}{|B(\bar{w}, w)|} h(w) d\mu(w) &= \int_{\mathbb{C}_+} |B(\bar{s}, \bar{w})|^2 |B(\bar{w}, w)|^t |B(\bar{w}, w)|^{-1} d\mu(w) \\ &= \int_{\mathbb{C}_+} |B(\bar{s}, \bar{w})|^2 |B(\bar{w}, w)|^{t-1} d\mu(w) \\ &\leq C B(\bar{s}, s)^{2+t-1-1} = C B(\bar{s}, s)^t = Ch(s), \end{aligned} \quad (2)$$

for some constant $C > 0$. From Schur's theorem [7], it follows that Q_1 is a bounded operator on $L^2(\mathbb{C}_+, d\mu)$. Moreover (1) and (2) also yield the boundedness of \mathcal{V}_1 . \square

The boundedness of Q_1 or \mathcal{V}_1 on $L^2(\mathbb{C}_+, d\mu)$ enables us to use Fubini's theorem [5]. Let $\phi, g \in L^2(\mathbb{C}_+, d\mu)$. Then

$$\begin{aligned} \langle \mathcal{V}_1 \phi, g \rangle &= \int_{\mathbb{C}_+} \left(3 \int_{\mathbb{C}_+} \frac{|B(\bar{s}, \bar{w})|^2}{|B(\bar{w}, w)|} \phi(s) d\mu(s) \right) \overline{g(w)} d\mu(w) \\ &= \int_{\mathbb{C}_+} \left(3 \int_{\mathbb{C}_+} \frac{|B(\bar{w}, \bar{s})|^2}{|B(\bar{w}, w)|} g(w) d\mu(w) \right) \phi(s) d\mu(s) \\ &= \langle \phi, Q_1 g \rangle, \end{aligned} \quad (3)$$

where the second equality of (3) follows from Fubini's theorem because

$$\begin{aligned} 3 \int_{\mathbb{C}_+} \int_{\mathbb{C}_+} \left| \frac{B(\bar{w}, \bar{s})^2}{B(\bar{w}, w)} \phi(s) g(w) \right| d\mu(s) d\mu(w) \\ \leq \|Q_1\| \|g\| \|\phi\| < \infty. \end{aligned}$$

Therefore, the adjoint operator of \mathcal{V}_1 on $L^2(\mathbb{C}_+, d\mu)$ is equal to Q_1 .

Lemma 5. For $\phi \in L^2(\mathbb{C}_+, d\mu)$,

$$\int_{\mathbb{C}_+} f(w)\overline{\phi(w)}d\mu(w) = \int_{\mathbb{C}_+} f(w)\overline{\mathcal{V}_1\phi(w)}d\mu(w)$$

for all $f \in L^2_a(\mathbb{C}_+)$.

Proof. As $Q_1f = f$ for $f \in L^2_a(\mathbb{C}_+)$, we have

$$\int_{\mathbb{C}_+} f(w)\overline{\phi(w)}d\mu(w) = \langle Q_1f, \phi \rangle = \langle f, \mathcal{V}_1\phi \rangle = \int_{\mathbb{C}_+} f(w)\overline{\mathcal{V}_1\phi(w)}d\mu(w).$$

□

5 Little Hankel operators

In this section, we establish that if $\phi \in L^2(\mathbb{C}_+, d\mu)$, then the little Hankel operator $h_{\overline{\phi}}$ is bounded if and only if $(\mathcal{V}_1\phi)(w)$ is bounded in \mathbb{C}_+ . Let $H^\infty(\mathbb{C}_+)$ be the space of bounded analytic functions on \mathbb{C}_+ . It is not difficult to verify that $H^\infty(\mathbb{C}_+) = WH^\infty(\mathbb{D})$ and $H^\infty(\mathbb{C}_+)$ is dense in $L^2_a(\mathbb{C}_+)$.

Proposition 3. If $\phi \in L^2(\mathbb{C}_+, d\mu)$, then $h_{\overline{\phi}} = h_{\overline{P_+\phi}}$ in the sense that $h_{\overline{\phi}}g = h_{\overline{P_+\phi}}g$ for all $g \in H^\infty(\mathbb{C}_+)$.

Proof. Let $h \in L^2_a(\mathbb{C}_+)$ and $g \in H^\infty(\mathbb{C}_+)$. Then

$$\begin{aligned} \langle h_{\overline{\phi}}g, \overline{h} \rangle &= \langle \overline{P_+(\phi g)}, \overline{h} \rangle = \langle \overline{\phi g}, \overline{h} \rangle = \langle gh, \phi \rangle \\ &= \langle gh, P_+\phi \rangle = \langle \overline{P_+\phi g}, \overline{h} \rangle = \langle \overline{P_+\phi g}, \overline{P_+h} \rangle \\ &= \langle \overline{P_+(P_+\phi g)}, \overline{h} \rangle = \langle h_{\overline{P_+\phi}}g, \overline{h} \rangle. \end{aligned}$$

Hence $h_{\overline{\phi}}g = h_{\overline{P_+\phi}}g$ for all $g \in H^\infty(\mathbb{C}_+)$. □

Lemma 6. Let $G(s) \in L^\infty(\mathbb{C}_+)$. Then the little Hankel operator Γ_G determined on $L^2_a(\mathbb{C}_+)$ by G is equivalent to the little Hankel operator Γ_ϕ determined on $L^2_a(\mathbb{D})$ by the function $\phi(z) = \left(\frac{1+\overline{z}}{1+z}\right)^2 G(Mz)$.

Proof. Notice that the sequence of vectors $\{\sqrt{n+1}z^n\}_{n=0}^\infty$ forms an orthonormal basis for $L^2_a(\mathbb{D})$. Then

$$\begin{aligned} \Gamma_G(W(\sqrt{n+1}z^n)) &= P_+ \left(G\mathcal{J} \left(\frac{2}{\sqrt{\pi}} \left(\frac{1-s}{1+s} \right)^n \frac{1}{(1+s)^2} \sqrt{n+1} \right) \right) \\ &= WPW^{-1} \left(G(s) \frac{2}{\sqrt{\pi}} \left(\frac{1-\overline{s}}{1+\overline{s}} \right)^n \frac{1}{(1+\overline{s})^2} \sqrt{n+1} \right) \\ &= W\Gamma_{\left(\frac{1+\overline{z}}{1+z}\right)^2 G(Mz)}(\sqrt{n+1}z^n) \text{ for all } n \geq 0. \end{aligned}$$

Thus Γ_G is unitarily equivalent to Γ_ϕ where $\phi(z) = \left(\frac{1+\bar{z}}{1+z}\right)^2 G(Mz)$. The result follows. \square

Proposition 4. *If $\phi \in L^\infty(\mathbb{C}_+)$, then $\bar{h}_\phi W = Wh_{\bar{\phi} \circ M}$.*

Proof. For $\phi \in L^\infty(\mathbb{C}_+)$, notice that $\bar{h}_\phi = \mathcal{J}\mathcal{S}_\phi$ and $\Gamma_\phi = \mathcal{S}_{\mathcal{J}\phi}$, where \mathcal{J} is the mapping from $L^2(\mathbb{C}_+, d\mu)$ into itself defined by $\mathcal{J}f(s) = f(\bar{s})$. Then from Lemma (6), we obtain

$$W^{-1}\mathcal{J}\bar{h}_{\mathcal{J}\phi}W = Jh_{J\left(\left(\frac{1+\bar{z}}{1+z}\right)^2(\bar{\phi} \circ M)(z)\right)}.$$

Hence

$$(W^{-1}\mathcal{J}W)(W^{-1}\bar{h}_{\mathcal{J}\phi}W) = Jh_{J\left(\left(\frac{1+\bar{z}}{1+z}\right)^2(\bar{\phi} \circ M)(z)\right)}.$$

Thus

$$J[J(W^{-1}\bar{h}_{\mathcal{J}\phi}W)] = J\left(Jh_{J\left(\left(\frac{1+\bar{z}}{1+z}\right)^2(\bar{\phi} \circ M)(z)\right)}\right).$$

Therefore

$$W^{-1}\bar{h}_{\mathcal{J}\phi}W = h_{J\left(\left(\frac{1+\bar{z}}{1+z}\right)^2(\bar{\phi} \circ M)(z)\right)}.$$

Hence

$$\bar{h}_{\mathcal{J}\phi}W = Wh_{J(u(\bar{\phi} \circ M))}, \quad (4)$$

where $u(z) = \left(\frac{1+\bar{z}}{1+z}\right)^2 = J(M' \circ M)(z)M'(z)$. Now from (4), it follows that

$$\bar{h}_\phi W = Wh_{J(u(\mathcal{J}\bar{\phi} \circ M))}. \quad (5)$$

Now

$$Ju = J(J(M' \circ M)M') = (M' \circ M)JM'.$$

Hence

$$(Ju \circ M) = (M' \circ M \circ M)(JM' \circ M) = M'(J(M' \circ M)).$$

Thus

$$\begin{aligned} (Ju)(Ju \circ M) &= (M' \circ M)(JM')M'(J(M' \circ M)) \\ &= (M' \circ M)M'J[(M' \circ M)M'] = 1. \end{aligned} \quad (6)$$

Further notice that

$$W^{-1}\bar{\phi} = (-1)\sqrt{\pi}(\bar{\phi} \circ M)M'.$$

Hence

$$J(W^{-1}\bar{\phi}) = (-1)\sqrt{\pi}(J\bar{\phi} \circ M)(JM').$$

This implies

$$WJW^{-1}\bar{\phi} = (-1)\sqrt{\pi}\frac{(-1)}{\sqrt{\pi}}(J\bar{\phi})(JM' \circ M)M' = (J\bar{\phi})(J(M' \circ M))M'.$$

Thus

$$\mathcal{J}\bar{\phi} = WJW^{-1}\bar{\phi} = u(J\bar{\phi}).$$

Hence

$$(\mathcal{J}\bar{\phi}) \circ M = (u \circ M)(J\bar{\phi} \circ M) = (u \circ M)J(\bar{\phi} \circ M).$$

Therefore

$$\begin{aligned} J((\mathcal{J}\bar{\phi}) \circ M) &= (J(u \circ M))(JJ(\bar{\phi} \circ M)) \\ &= (J(u \circ M))(\bar{\phi} \circ M) \\ &= ((Ju) \circ M)(\bar{\phi} \circ M). \end{aligned}$$

Form (5), we obtain

$$\begin{aligned} \hbar_{\bar{\phi}}W &= Wh_{J(u(\mathcal{J}\bar{\phi} \circ M))} = Wh_{(Ju)(J(\mathcal{J}\bar{\phi} \circ M))} \\ &= Wh_{Ju[(Ju \circ M)(\bar{\phi} \circ M)]} = Wh_{[(Ju)(Ju \circ M)](\bar{\phi} \circ M)}. \end{aligned}$$

From (6), it follows that $\hbar_{\bar{\phi}}W = Wh_{\bar{\phi} \circ M}$. \square

For $\phi \in L^2(\mathbb{C}_+, d\mu)$, it is not difficult to show that $(\mathcal{V}_1\phi)(w) = 3\langle \bar{b}_w, \hbar_{\bar{\phi}}b_w \rangle$. For $z \in \mathbb{D}$, $f \in L^2(\mathbb{D}, dA)$, define

$$(Vf)(z) = 3(1 - |z|^2)^2 \int_{\mathbb{D}} \frac{f(w)}{(1 - z\bar{w})^4} dA(w).$$

Proposition 5. *Let $\phi \in L^2(\mathbb{C}_+, d\mu)$, then $(\mathcal{V}_1\phi)(w) = V(\phi \circ M)(a)$, for all $a \in \mathbb{D}$.*

Proof. Let $\phi \in L^2(\mathbb{C}_+, d\mu)$ and $w = M\bar{a}$, $a \in \mathbb{D}$, $w \in \mathbb{C}_+$. Then

$$\begin{aligned} \mathcal{V}_1\phi(w) &= 3\langle \bar{b}_w, \hbar_{\bar{\phi}}b_w \rangle = 3\langle \overline{Wk_a}, \hbar_{\bar{\phi}}Wk_a \rangle = 3\langle W\bar{k}_a, \hbar_{\bar{\phi}}Wk_a \rangle \\ &= 3\langle \bar{k}_a, W^{-1}\hbar_{\bar{\phi}}Wk_a \rangle = 3\langle \bar{k}_a, \hbar_{\bar{\phi} \circ M}k_a \rangle = 3\langle \bar{k}_a, \overline{h_{\bar{\phi} \circ M}k_a} \rangle \\ &= V(\phi \circ M)(a), \end{aligned}$$

for all $a \in \mathbb{D}$. \square

Proposition 6. For $\phi \in L^2(\mathbb{C}_+, d\mu)$,

(i) $\mathcal{V}_1 P_+ = \mathcal{V}_1$.

(ii) $P_+ \mathcal{V}_1 = P_+$.

(iii) $\mathcal{V}_1^2 = \mathcal{V}_1$.

Proof. From Proposition 3, we obtain

$$\mathcal{V}_1 P_+ \phi = 3 \langle \bar{b}_w, \bar{h}_{P_+ \phi} b_w \rangle = 3 \langle \bar{b}_w, \bar{h}_\phi b_w \rangle = \mathcal{V}_1 \phi,$$

for $\phi \in L^2(\mathbb{C}_+, d\mu)$. This proves (i). To prove (ii), let $\phi, g \in L^2(\mathbb{C}_+, d\mu)$ and $g = g_1 + g_2$ where $g_1 \in L^2_a(\mathbb{C}_+)$ and $g_2 \in (L^2_a(\mathbb{C}_+))^\perp$. Then

$$\begin{aligned} \langle P_+ \mathcal{V}_1 \phi, g \rangle &= \langle \mathcal{V}_1 \phi, P_+ g \rangle = \langle \mathcal{V}_1 \phi, g_1 \rangle = \int_{\mathbb{C}_+} (\mathcal{V}_1 \phi)(w) \overline{g_1(w)} d\mu(w) \\ &= \pi \int_{\mathbb{D}} [(\mathcal{V}_1 \phi) \circ M](z) \overline{(g_1 \circ M)(z)} |M'(z)|^2 dA(z) \\ &= \pi \int_{\mathbb{D}} [V(\phi \circ M)](z) \overline{(g_1 \circ M)(z)} |M'(z)|^2 dA(z). \end{aligned}$$

Under the complex integral pairing with respect to dA , we have $V = P_2^*$ where $P_2 h(z) = 3 \int_{\mathbb{D}} \frac{(1 - |u|^2)^2}{(1 - z\bar{u})^4} h(u) dA(u)$ is a projection from $L^1(\mathbb{D}, dA)$ onto $L^1_a(\mathbb{D})$. From Fubini's theorem [5] and the fact that both P and P_2 reproduce analytic functions it follows that $PV = P$, where P is the Bergman projection from $L^2(\mathbb{D}, dA)$ onto $L^2_a(\mathbb{D})$. Thus for $\phi, g \in L^2(\mathbb{C}_+, d\mu)$,

$$\begin{aligned} \langle P_+ \mathcal{V}_1 \phi, g \rangle &= \pi \int_{\mathbb{D}} [V(\phi \circ M)](z) \overline{(g_1 \circ M)(z)} |M'(z)|^2 dA(z) \\ &= \pi \int_{\mathbb{D}} V[(\phi \circ M)M'](z) \overline{(g_1 \circ M)(z)M'(z)} dA(z) \\ &= \int_{\mathbb{D}} V[(-1)\sqrt{\pi}(\phi \circ M)M'](z) \overline{(-1)\sqrt{\pi}(g_1 \circ M)(z)M'(z)} dA(z) \\ &= \int_{\mathbb{D}} V(W^{-1}\phi)(z) \overline{(W^{-1}g_1)(z)} dA(z) \\ &= \langle VW^{-1}\phi, W^{-1}g_1 \rangle = \langle VW^{-1}\phi, W^{-1}P_+g_1 \rangle = \langle VW^{-1}\phi, PW^{-1}g_1 \rangle \\ &= \langle PVW^{-1}\phi, W^{-1}g_1 \rangle = \langle PW^{-1}\phi, W^{-1}g_1 \rangle = \langle WPW^{-1}\phi, g_1 \rangle \\ &= \langle P_+ \phi, g_1 \rangle = \langle P_+^2 \phi, g_1 \rangle = \langle P_+ \phi, P_+ g_1 \rangle \\ &= \langle P_+ \phi, P_+ g \rangle = \langle P_+^2 \phi, g \rangle = \langle P_+ \phi, g \rangle. \end{aligned}$$

Thus $P_+ \mathcal{V}_1 \phi = P_+ \phi$ for all $\phi \in L^2(\mathbb{C}_+, d\mu)$ and therefore $P_+ \mathcal{V}_1 = P_+$. This proves (ii). To prove (iii), notice that

$$\begin{aligned} (\mathcal{V}_1^2 \phi)(w) &= \mathcal{V}_1(\mathcal{V}_1 \phi)(w) = 3\langle \bar{b}_w, \bar{h}_{\mathcal{V}_1 \phi} b_w \rangle = 3\langle \bar{b}_w, \bar{h}_{P_+ \mathcal{V}_1 \phi} b_w \rangle \\ &= 3\langle \bar{b}_w, \bar{h}_{P_+ \phi} b_w \rangle = 3\langle \bar{b}_w, \bar{h}_\phi b_w \rangle = (\mathcal{V}_1 \phi)(w) \end{aligned}$$

for all $w \in \mathbb{C}_+$ and $\phi \in L^2(\mathbb{C}_+, d\mu)$. Hence $\mathcal{V}_1^2 = \mathcal{V}_1$. □

Proposition 7. *Let $a \in \mathbb{D}, \bar{f} \in \overline{L_a^2(\mathbb{D})}$ and $f = W^{-1}g, g \in L_a^2(\mathbb{C}_+)$. Then*

$$h_{\phi \circ M}^* \bar{f}(a) = c_a \langle \bar{h}_\phi \bar{g}, B_w \rangle,$$

for all $g \in L_a^2(\mathbb{C}_+)$ and for some constant c_a .

Proof. Let $a \in \mathbb{D}, \bar{f} \in \overline{L_a^2(\mathbb{D})}$ and $f = W^{-1}g, g \in L_a^2(\mathbb{C}_+)$. Then by Lemma 2, there exists a constant $\alpha, |\alpha| = 1$ such that

$$\begin{aligned} h_\phi^* \bar{f}(a) &= \langle h_\phi^* \bar{f}, K_a \rangle = \langle \bar{f}, h_\phi K_a \rangle = \langle W \bar{f}, W h_\phi K_a \rangle \\ &= \|K_a\| \langle W \bar{f}, W h_\phi k_a \rangle = \|K_a\| \langle \bar{g}, W h_\phi W^{-1} b_w \rangle = \|K_a\| \langle \bar{g}, \bar{h}_{\phi \circ M} b_w \rangle \\ &= \|K_a\| \langle \bar{h}_{\phi \circ M}^* \bar{g}, b_w \rangle = \alpha \|K_a\| \left\langle \bar{h}_{\phi \circ M}^* \bar{g}, \frac{B_w}{\|B_w\|} \right\rangle = \frac{\alpha \|K_a\|}{\|B_w\|} \langle \bar{h}_{\phi \circ M}^* \bar{g}, B_w \rangle \\ &= c_a \langle \bar{h}_{\phi \circ M}^* \bar{g}, B_w \rangle, \end{aligned}$$

where $c_a = \frac{\alpha \|K_a\|}{\|B_w\|}$. Thus,

$$h_{\phi \circ M}^* \bar{f}(a) = c_a \langle \bar{h}_\phi \bar{g}, B_w \rangle.$$

□

Theorem 1. *Suppose $\phi \in L^2(\mathbb{C}_+, d\mu)$. Then \bar{h}_ϕ is bounded if and only if $(\mathcal{V}_1 \phi)(w)$ is bounded in \mathbb{C}_+ and there is a constant $C > 0$ such that $C^{-1} \|\mathcal{V}_1 \phi\|_\infty \leq \|\bar{h}_\phi\| \leq C \|\mathcal{V}_1 \phi\|_\infty$.*

Proof. Notice that $b_w \in L^2(\mathbb{C}_+, d\mu)$ and $\|b_w\|_2 = 1$. Hence $|(\mathcal{V}_1 \phi)(w)| = 3|\langle \bar{b}_w, \bar{h}_\phi b_w \rangle| \leq 3\|b_w\|_2 \|\bar{h}_\phi\| \|b_w\|_2 = 3\|b_w\|_2^2 \|\bar{h}_\phi\| = 3\|\bar{h}_\phi\|$. Further, $\bar{h}_\phi = \bar{h}_{P_+ \phi} = \bar{h}_{P_+ \mathcal{V}_1 \phi} = \bar{h}_{\mathcal{V}_1 \phi}$. Thus $\mathcal{V}_1 \phi \in L^\infty(\mathbb{C}_+)$ implies that \bar{h}_ϕ is bounded with $\|\bar{h}_\phi\| \leq \|\mathcal{V}_1 \phi\|_\infty$. The result follows since $\bar{h}_\phi = \bar{h}_{\mathcal{V}_1 \phi}$ for all $\phi \in L^2(\mathbb{C}_+, d\mu)$. □

Theorem 2. *Suppose $\phi \in L^2(\mathbb{C}_+, d\mu)$ such that \bar{h}_ϕ is bounded. Then*

$$\frac{1}{3} \limsup_{Re w \rightarrow 0} |\mathcal{V}_1 \phi(w)| \leq \|\bar{h}_\phi\|_e.$$

Proof. Consider a compact operator T from $L_a^2(\mathbb{C}_+, d\mu)$ to $\overline{L_a^2(\mathbb{C}_+, d\mu)}$ arbitrarily. Since $b_{\bar{w}} \rightarrow 0$ weakly in $L_a^2(\mathbb{C}_+)$ as $Re w \rightarrow 0$, we obtain $\|Tb_{\bar{w}}\| \rightarrow 0$ as $Re w \rightarrow 0$. Hence,

$$\|\tilde{h}_{\bar{\phi}} - T\| \geq \limsup_{Re w \rightarrow 0} \|(\tilde{h}_{\bar{\phi}} - T)b_{\bar{w}}\| \geq \limsup_{Re w \rightarrow 0} \|\tilde{h}_{\bar{\phi}}b_{\bar{w}}\|. \tag{7}$$

Since (7) holds for every compact operator T , it follows that,

$$\|\tilde{h}_{\bar{\phi}}\|_e \geq \limsup_{Re w \rightarrow 0} \|\tilde{h}_{\bar{\phi}}b_{\bar{w}}\|. \tag{8}$$

On the other hand,

$$|(\mathcal{V}_1\phi)(w)| = 3|\langle b_{\bar{w}}, \tilde{h}_{\bar{\phi}}b_{\bar{w}} \rangle| \leq 3\|\tilde{h}_{\bar{\phi}}b_{\bar{w}}\|. \tag{9}$$

From (8) and (9), the theorem follows. □

6 Main Result

In this section, we give estimates for the essential norm of bounded little Hankel operators on the Bergman space $L_a^2(\mathbb{C}_+, d\mu)$ in terms of the function $\mathcal{V}_1\phi$ and applications of the result are also derived. Assume that $\tilde{h}_{\bar{\phi}}$ is bounded operator from $L_a^2(\mathbb{C}_+, d\mu)$ to $\overline{L_a^2(\mathbb{C}_+, d\mu)}$. The following holds:

Theorem 3. *Suppose $\phi \in L^2(\mathbb{C}_+, d\mu)$ and $\tilde{h}_{\bar{\phi}}$ is bounded. Then*

$$\|\tilde{h}_{\bar{\phi}}\|_e \leq C \limsup_{Re w \rightarrow 0} |\mathcal{V}_1\phi(w)|.$$

Proof. For $f \in \overline{L_a^2(\mathbb{D})}$ and $0 < r < 1$, define

$$F_r f(z) = \int_{\mathbb{D}} \left(\int_{r\mathbb{D}} \frac{1}{(1 - z\bar{u})^2} \frac{1}{(1 - v\bar{u})^2} V\phi(u) dA(u) \right) f(v) dA(v).$$

Then PF_r is a compact operator from $\overline{L_a^2(\mathbb{D})}$ to $L_a^2(\mathbb{D})$ because

$$\begin{aligned} & \int_{\mathbb{D}} \int_{\mathbb{D}} \left| \int_{r\mathbb{D}} \frac{1}{(1 - z\bar{u})^2} \frac{1}{(1 - v\bar{u})^2} V\phi(u) dA(u) \right|^2 dA(z) dA(v) \\ & \leq \|V\phi\|_{\infty}^2 \int_{r\mathbb{D}} \int_{\mathbb{D}} \int_{\mathbb{D}} \frac{1}{|1 - z\bar{u}|^4} \frac{1}{|1 - v\bar{u}|^4} dA(z) dA(v) dA(u) \end{aligned}$$

$$= \|V\phi\|_\infty^2 \int_{r\mathbb{D}} \frac{1}{(1-|u|^2)^2} dA(u) < \infty$$

and the result follows from Theorem 3.5 in [7]. Thus, we have

$$\|h_{\bar{\phi}}\|_e = \|h_{\bar{\phi}}^*\|_e = \|h_{\bar{\phi}}^* - PF_r\|_e \leq \|h_{\bar{\phi}}^* - PF_r\|.$$

Moreover, $Ph_{\bar{\phi}}^* = h_{\bar{\phi}}^*$ yields,

$$\begin{aligned} \|h_{\bar{\phi}}^* - PF_r\| &= \sup_{f \in L_a^2(\mathbb{D})} \frac{\|Ph_{\bar{\phi}}^*f - PF_rf\|}{\|f\|} \\ &\leq \sup_{f \in L_a^2(\mathbb{D})} \frac{\|h_{\bar{\phi}}^*f - F_rf\|}{\|f\|}. \end{aligned}$$

Define

$$K_r(z, v) = \int_{\mathbb{D}/r\mathbb{D}} \frac{1}{(1-z\bar{u})^2} \frac{1}{(1-v\bar{u})^2} V\phi(u) dA(u),$$

and $K_r^+(z, v) = |K_r(z, v)|$. Let G_r (respectively G_r^+) be the integral operator on $L^2(\mathbb{D}, dA)$ with kernel K_r (respectively K_r^+). Then $G_rf = h_{\bar{\phi}}^*f - F_rf$ for any $f \in \overline{L_a^2(\mathbb{D})}$. For details see [7]. Thus

$$\|h_{\bar{\phi}}\|_e \leq \|G_r^+\|.$$

Using Schur's theorem, we will obtain the operator norm of G_r^+ on $L^2(\mathbb{D}, dA)$. Take $t > 0$ and $h(z) = \frac{1}{(1-|z|^2)^{2t}}$. Then we have,

$$\begin{aligned} &\int_{\mathbb{D}} K_r^+(z, v) h(v) dA(v) \\ &= \int_{\mathbb{D}} \left| \int_{\mathbb{D}/r\mathbb{D}} \frac{1}{(1-z\bar{u})^2} \frac{1}{(1-v\bar{u})^2} V\phi(u) dA(u) \right| \frac{1}{(1-|v|^2)^{2t}} dA(v) \\ &\leq \left(\sup_{r < |u| < 1} |V\phi(u)| \right) \int_{\mathbb{D}} \int_{\mathbb{D}/r\mathbb{D}} \left| \frac{1}{(1-z\bar{u})^2} \frac{1}{(1-v\bar{u})^2} \right| \frac{1}{(1-|v|^2)^{2t}} dA(u) dA(v). \end{aligned}$$

Since

$$\begin{aligned} &\int_{\mathbb{D}} \int_{\mathbb{D}/r\mathbb{D}} \left| \frac{1}{(1-z\bar{u})^2} \frac{1}{(1-v\bar{u})^2} \right| \frac{1}{(1-|v|^2)^{2t}} dA(u) dA(v) \\ &\leq \int_{\mathbb{D}} \int_{\mathbb{D}} \frac{1}{|1-v\bar{u}|^2} \frac{1}{(1-|v|^2)^{2t}} \frac{1}{|1-z\bar{u}|^2} dA(v) dA(u) \\ &\leq Ch(z), \end{aligned}$$

we obtain from [3] that,

$$\int_{\mathbb{D}} K_r^+(z, v)h(v)dA(v) \leq C \left(\sup_{r<|u|<1} |V\phi(u)| \right) h(z).$$

Thus using Schur’s theorem [7], we have

$$\|G_r^+\| \leq C \sup_{r<|u|<1} |V\phi(u)|.$$

Thus

$$\|h_{\bar{\phi}}\|_e \leq C \sup_{r<|u|<1} |V\phi(u)|,$$

for any $0 < r < 1$. Letting $r \rightarrow 1$, we obtain

$$\|h_{\bar{\phi}}\|_e \leq C \limsup_{u \rightarrow \partial\mathbb{D}} |V\phi(u)|.$$

Hence

$$\begin{aligned} \|h_{\bar{\phi}}\|_e &= \inf\{\|h_{\bar{\phi}} - T\| : T \text{ is compact}\} \\ &= \inf\{\|W^{-1}h_{\bar{\phi}}W - W^{-1}TW\| : T \text{ is compact}\} \\ &= \{\|h_{\bar{\phi} \circ M}W - L\| : L \text{ is compact in } \mathcal{L}(L_a^2(\mathbb{D}))\} \\ &= \|h_{\bar{\phi} \circ M}\|_e \\ &\leq C \limsup_{a \rightarrow \partial\mathbb{D}} |V(\phi \circ M)(a)| = C \limsup_{Re w \rightarrow 0} |(\mathcal{V}_1\phi)(w)|. \end{aligned}$$

□

Corollary 1. *Let $\phi \in L^2(\mathbb{C}_+, d\mu)$. Then $h_{\bar{\phi}}$ is a compact operator from $L_a^2(\mathbb{C}_+, d\mu)$ to $\overline{L_a^2(\mathbb{C}_+, d\mu)}$ if and only if $\mathcal{V}_1\phi(w) \rightarrow 0$ as $Re w \rightarrow 0$.*

Proof. Suppose $h_{\bar{\phi}}$ is compact. Since $h_{\bar{\phi}}$ is bounded and $\|h_{\bar{\phi}}\|_e = 0$. It thus follows from Theorem 2, that $\limsup_{Re w \rightarrow 0} |\mathcal{V}_1\phi(w)| = 0$. That is, $\mathcal{V}_1\phi(w) \rightarrow 0$ as $Re w \rightarrow 0$. On the other hand, $\mathcal{V}_1\phi(w) \rightarrow 0$ and since $\mathcal{V}_1\phi$ is a continuous functions, we obtain that $\mathcal{V}_1\phi$ is bounded. Therefore, $h_{\bar{\phi}}$ is bounded. Hence, from Theorem 3, we obtain $\|h_{\bar{\phi}}\|_e = 0$. Thus $h_{\bar{\phi}}$ is compact. □

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